Charge dispersion MPGD Readout & ILC-TPC Prototype Test Beam Studies

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Motivation & overview

- ILC tracker goal ∆(1/p_T) ≤ 5.10⁻⁵ (GeV/c)⁻¹
 => MPGD-TPC ∆(1/p_T) ≤ 1.5 x 10⁻⁴ (GeV/c)⁻¹
- TDR TPC: 200 pads; $\sigma_{Tr} \sim 100 \ \mu m$ ($\approx 2 \ m \ drift$), pad size 2 x 6 mm² => Total TPC pad count ~1.5 x 10⁶
- R&D shows 2 mm too wide for 100 μm resolution with normal readout. Ways to improve the MPGD-TPC resolution:
- Under consideration narrower 1 mm x 6 mm pads (3 x 10⁶ total). R&D issues: High density electronics, increased heat load, TPC endcap mass etc.
- 2) Alternative: Disperse avalanche charge to improve resolution for wide pads. Development of TPC readout with charge dispersion in MPGDs with a resistive anode.
 - Charge dispersion demonstrated in cosmic ray TPC tests with no magnet.
 - The new TPC readout concept was tested in a beam test last October. 1 T superconducting magnet & 4 GeV/c hadron test beam at KEK PS.
 - Two TPCs: Multi Technology Test TPC MT3 TPC (MPI Munich) + Carleton TPC with Micromegas (Saclay) & GEMs(Saga University).
 - Two weeks of beam data in October2005.
- 3) Magnetic field performance of MPGD-TPC with charge dispersion readout in a test beam.

Diffusion sets the limit on TPC momentum resolution

•The physics limit of TPC resolution comes from transverse diffusion: $D^2 + z$

$$\sigma_x^2 \approx \frac{D_{Tr}^2 \cdot z}{N_{eff}} N_{eff}$$
 = effective electron statistics.

•For best resolution, choose a gas with smallest diffusion

The rule applies to the wire TPCs which use induced cathode pad signals for position determination. But **ExB** & track angle systematic effects degrade wire TPC resolution.

ExB effect does not limit the MPGD-TPC. But there are no comparable induced cathode pad signals.

The MPGD-TPC resolution is limited by pad width *w*. The resolution gets worse for wide pads in absence of diffusion.

$$\sigma_x^2 \Rightarrow \frac{w^2}{12} \text{ as } z \Rightarrow 0$$

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Pad width limits the MPGD-TPC resolution ExB angle effects limit the wire/pad TPC resolution



Direct signal on the MPGD anode pad For small diffusion, less precise centroid for wide pads Induced cathode signal determined by geometry Accurate centroid determination possible with wide pads

$$\sigma_x^2 \approx \sigma_0^2 + \frac{1}{N_{eff}} \left[D_{Tr}^2 z + w^2 / 12 \right]$$

$$\sigma_x^2 \approx \sigma_0^2 + \frac{D_{Tr}^2 \cdot z}{N_{eff}}$$

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Improving MPGD TPC resolution without resorting to narrower pads

- •Disperse track charge after gas gain to improve centroid determination with wide pads.
- •For the GEM, large transverse diffusion in the high E-field field in transfer and induction gaps provides a natural mechanism to disperse the cluster charge.
- •Measurements with prototype GEM-TPCs indicate that ~ 1 mm wide pads may be needed for TPC operation in high magnetic fields.
- •Explore other concepts to disperse the charge

Charge dispersion - a geometrical pad signal induction mechanism making position sensing insensitive to pad width.

Position sensing from charge dispersion in a MPGD with a resistive anode

Position sensing on a resistive anode proportional wire from charge division

Telegraph equation (1-D):

Deposit point charge at t=0

Solution for charge density $(L \sim 0)$

$$\frac{L}{R}\frac{\partial^2 Q}{\partial t^2} + \frac{\partial Q}{\partial t} = \frac{1}{RC}\frac{\partial^2 Q}{\partial x^2}$$
$$Q(x,t) = \sqrt{\frac{RC}{4\pi t}}e^{\frac{-x^2 RC}{4t}}$$

Generalize charge division on a resistive wire to 2 D

Position sensing from charge dispersion in MPGDs with a resistive anode

Equivalent to Telegraph equation in 2-D

$$\frac{\partial Q}{\partial t} = \frac{1}{RC} \left[\frac{\partial^2 Q}{\partial r^2} + \frac{1}{r} \frac{\partial Q}{\partial r} \right]$$

Solution for charge density in 2-D

$$Q(r,t) = \frac{RC}{2t}e^{\frac{-r^2RC}{4t}}$$

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Charge dispersion in a MPGD with a resistive anode

•Modified MPGD anode with a high resistivity film bonded to a readout plane with an insulating spacer.

•2-dimensional continuous RC network defined by material properties & geometry.

•Point charge at r = 0 & t = 0 disperses with time.

•Time dependent anode charge density sampled by readout pads.

Equation for surface charge density function on the 2-dim. continuous RC network:

$$\frac{\partial \rho}{\partial t} = \frac{1}{RC} \left[\frac{\partial^2 \rho}{\partial r^2} + \frac{1}{r} \frac{\partial \rho}{\partial r} \right]$$
$$\Rightarrow \rho(r,t) = \frac{RC}{2t} e^{\frac{-r^2 RC}{4t}}$$

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M.S.Dixit et.al., Nucl. Instrum. Methods A518 (2004) 721.

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Simulation of charge dispersion phenomenon

- •The charge dispersion equation describe the time evolution of a point like charge deposited on the MPGD resistive anode at t = 0.
- •To compare to experiment, one needs to include the effects of:
 - •Longitudinal & transverse diffusion in the gas.
 - •Intrinsic rise time T_{rise} of the detector charge pulse.
 - •The effect of preamplifier rise and fall times $t_r \& t_f$.
 - •And for particle tracks, the effects of primary ionization clustering.

The simulation for a single charge cluster

The charge density function for a point charge in Cartesian coordinates:

$$\rho_{\delta}(x, y, t) = \frac{\tau}{4\pi t} \exp\left[-\tau \left(x^2 + y^2\right)/4t\right] \text{ where } \tau = RC$$

Physics effects included in simulation in two parts: 1) as effects which depend on spatial coordinates x & y, or; 2) as effects which depend on time.
1) The spatial effects function includes charge dispersion phenomena & transverse size *w* of the charge cluster due to transverse diffusion.

 $Q_{pad}(t)$ is the pad signal from charge dispersion when a charge Nq_e of size *w* is deposited on the anode at t = 0;

$$Q_{pad} = \frac{Nq_e}{4} \left[erf(\frac{x_{high}}{\sqrt{2}\sigma_{xy}}) - erf(\frac{x_{low}}{\sqrt{2}\sigma_{xy}}) \right] \left[erf(\frac{y_{high}}{\sqrt{2}\sigma_{xy}}) - erf(\frac{y_{low}}{\sqrt{2}\sigma_{xy}}) \right]$$
(1)

 $x_{high}, x_{low}, y_{high}, y_{low}$ define the pad boundaries & $\sigma_{xy} = \sqrt{2t} / \tau + w^2$

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$$I(t) = \frac{1}{2T_{rise}} \begin{bmatrix} exp(\sigma^2 a^2/2 - at) \left[erf\left(\frac{t - T_{rise} - \sigma^2 a}{\sigma\sqrt{2}}\right) + 1 \right] - \\ exp(\sigma^2 b^2/2 - bt) \left[erf\left(\frac{t - T_{rise} - \sigma^2 b}{\sigma\sqrt{2}}\right) + 1 \right] + \\ exp(\sigma^2 a^2/2 - a(t - T_{rise})) \left[erf\left(\frac{t - 2T_{rise} - \sigma^2 a}{\sigma\sqrt{2}}\right) + 1 \right] - \\ exp(\sigma^2 b^2/2 - b(t - T_{rise})) \left[erf\left(\frac{t - 2T_{rise} - \sigma^2 b}{\sigma\sqrt{2}}\right) + 1 \right] + \\ exp(\sigma^2 a^2/2 - at) \left[erf\left(\frac{t - \sigma^2 a}{\sigma\sqrt{2}}\right) + erf\left(\frac{t - T_{rise} - \sigma^2 a}{\sigma\sqrt{2}}\right) \right] - \\ exp(\sigma^2 b^2/2 - bt) \left[erf\left(\frac{t - \sigma^2 b}{\sigma\sqrt{2}}\right) + erf\left(\frac{t - T_{rise} - \sigma^2 b}{\sigma\sqrt{2}}\right) \right] - \\ exp(\sigma^2 b^2/2 - bt) \left[erf\left(\frac{t - \sigma^2 b}{\sigma\sqrt{2}}\right) + erf\left(\frac{t - T_{rise} - \sigma^2 b}{\sigma\sqrt{2}}\right) \right] - \\ exp(\sigma^2 b^2/2 - bt) \left[erf\left(\frac{t - \sigma^2 b}{\sigma\sqrt{2}}\right) + erf\left(\frac{t - T_{rise} - \sigma^2 b}{\sigma\sqrt{2}}\right) \right] - \\ exp(\sigma^2 b^2/2 - bt) \left[erf\left(\frac{t - \sigma^2 b}{\sigma\sqrt{2}}\right) + erf\left(\frac{t - T_{rise} - \sigma^2 b}{\sigma\sqrt{2}}\right) \right] - \\ exp(\sigma^2 b^2/2 - bt) \left[erf\left(\frac{t - \sigma^2 b}{\sigma\sqrt{2}}\right) + erf\left(\frac{t - T_{rise} - \sigma^2 b}{\sigma\sqrt{2}}\right) \right] - \\ exp(\sigma^2 b^2/2 - bt) \left[erf\left(\frac{t - \sigma^2 b}{\sigma\sqrt{2}}\right) + erf\left(\frac{t - T_{rise} - \sigma^2 b}{\sigma\sqrt{2}}\right) \right] - \\ exp(\sigma^2 b^2/2 - bt) \left[erf\left(\frac{t - \sigma^2 b}{\sigma\sqrt{2}}\right) + erf\left(\frac{t - T_{rise} - \sigma^2 b}{\sigma\sqrt{2}}\right) \right] - \\ exp(\sigma^2 b^2/2 - bt) \left[erf\left(\frac{t - \sigma^2 b}{\sigma\sqrt{2}}\right) + erf\left(\frac{t - T_{rise} - \sigma^2 b}{\sigma\sqrt{2}}\right) \right] - \\ exp(\sigma^2 b^2/2 - bt) \left[erf\left(\frac{t - \sigma^2 b}{\sigma\sqrt{2}}\right) + erf\left(\frac{t - T_{rise} - \sigma^2 b}{\sigma\sqrt{2}}\right) \right] - \\ exp(\sigma^2 b^2/2 - bt) \left[erf\left(\frac{t - \sigma^2 b}{\sigma\sqrt{2}}\right) + erf\left(\frac{t - T_{rise} - \sigma^2 b}{\sigma\sqrt{2}}\right) \right] - \\ exp(\sigma^2 b^2/2 - bt) \left[erf\left(\frac{t - \sigma^2 b}{\sigma\sqrt{2}}\right) + erf\left(\frac{t - T_{rise} - \sigma^2 b}{\sigma\sqrt{2}}\right) \right] - \\ exp(\sigma^2 b^2/2 - bt) \left[erf\left(\frac{t - \sigma^2 b}{\sigma\sqrt{2}}\right) + erf\left(\frac{t - T_{rise} - \sigma^2 b}{\sigma\sqrt{2}}\right) \right] - \\ exp(\sigma^2 b^2/2 - bt) \left[erf\left(\frac{t - \sigma^2 b}{\sigma\sqrt{2}}\right) + erf\left(\frac{t - T_{rise} - \sigma^2 b}{\sigma\sqrt{2}}\right) \right] - \\ exp(\sigma^2 b^2/2 - bt) \left[erf\left(\frac{t - \sigma^2 b}{\sigma\sqrt{2}}\right) + erf\left(\frac{t - T_{rise} - \sigma^2 b}{\sigma\sqrt{2}}\right) \right] + \\ exp(\sigma^2 b^2/2 - bt) \left[erf\left(\frac{t - \sigma^2 b}{\sigma\sqrt{2}}\right) + erf\left(\frac{t - T_{rise} - \sigma^2 b}{\sigma\sqrt{2}}\right) \right] + \\ exp(\sigma^2 b^2/2 - bt) \left[erf\left(\frac{t - \sigma^2 b}{\sigma\sqrt{2}}\right) + erf\left(\frac{t - T_{rise} - \sigma^2 b}{\sigma\sqrt{2}}\right) \right] + \\ exp(\sigma^2 b^2/2 - bt) \left[erf\left(\frac{t - \sigma^2 b}{\sigma\sqrt{2}}\right) + er$$

I(t) incorporates intrinsic rise time, longitudinal diffusion & electronics shaping times as time dependent effects. $a = 1/t_f$; $b = 1/t_f + 1/t_r$

(1) and (2) are convoluted numerically for the model simulation.

MPGD charge dispersion tests with a collimated point x ray source



•Point source ~ 50 μ m collimated 4.5 keV x rays. •Aleph TPC preamps. τ_{Rise} = 40 ns, τ_{Fall} = 2 μ s. •DAQ - 500 MHz Tektronix digital scope.

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Resistive anode Micromegas

530 k Ω/\Box Carbon loaded Kapton resistive anode was used with GEM. This was replaced with higher resistivity 1 M Ω/\Box Cermet for tests with Micromegas.



Charge dispersion signal for a GEM Simulation versus measurement (Ar+10%CO2) (2 x 6 mm² pads) Collimated ~ 50 μm 4.5 keV x-ray spot on pad centre.



Simulated primary pulse is normalized to the data.

Primary pulse normalization used for the simulated secondary pulse

GEM pad response function (PRF) for a point source Simulation compared to the measured PRF



lonization from 50 μm collimated x-ray spot.



Measured PRF deviates from simulation due to anode RC nonuniformities.

PRF - a measure of signal amplitude as a function of cluster position.

Double-GEM space point resolution with charge dispersion readout (Ar+10%CO2) Collimated ~ 4.5 keV x rays, Spot size ~ 50 μm



•GEM resolution ~ 70 μ m.

•Similar resolution measured for a Micromegas with a resistive anode readout using 2 mm x 6 mm pads

Cosmic ray TPC tests with charge dispersion

- 15 cm drift length with GEM or Micromegas readoutB=0
- •Ar+10% CO2 chosen to simulate low transverse diffusion in a magnetic field.

Aleph charge preamps. τ_{Rise}= 40 ns, τ_{Fall} = 2 μs.
200 MHz FADCs rebinned to digitization effectively at 25 MHz.
60 tracking pads (2 x 6 mm²) + 2 trigger pads (24 x 6 mm²).

The GEM-TPC resolution was first measured with conventional direct charge TPC readout.



The resolution was next measured with a charge dispersion resistive anode readout with a double-GEM & with a Micromegas endcap.

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GEM TPC charge dispersion simulation (B=0) Cosmic ray track, Z = 67 mm Ar+10%CO2



Centre pulse used for simulation normalization - no other free parameters.

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The pad response function (PRF) for a track

- The PRF is a measure of signal size as a function of track position relative to the pad.
- Unusual highly variable charge dispersion pulse shape; both the rise time & pulse amplitude depend on track position.
- We use pulse shape information to optimize the PRF.
- The PRF can, in principle, be determined from simulation.
- However, system RC nonuniformities & geometrical effects introduce bias in absolute position determination.
- The position bias can be corrected by calibration.
- PRF and bias determined empirically using a subset of data which was used for calibration. The remaining data used for resolution studies.

Track PRFs with GEM & Micromegas readout

- The PRFs are not Gaussian.
- The PRF depends on track position relative to the pad.
- PRF = PRF(x,z)
- PRF can be characterized by its FWHM $\Gamma(z)$ & base width $\Delta(z)$.
- PRFs determined from the data have been fitted to a functional form consisting of a ratio of two symmetric 4th order polynomials.

$$PRF[x,\Gamma(z),\Delta,a,b] = \frac{(1+a_2x^2+a_4x^4)}{(1+b_2x^2+b_4x^4)}$$

 $a_2 a_4 b_2 \& b_4$ can be written down in terms of Γ and $\Delta \&$ two scale parameters a & b.

GEM & Micromegas PRFs for tracks Ar+10%CO2 2x6 mm² pads

The pad response function amplitude for longer drift distances is lower due to Z dependent normalization.



Micromegas PRF is narrower due to the use of higher resistivity anode & smaller diffusion than GEM after avalanche gain

Track fit using the the PRF



Bias corrections for the GEM & for Micromegas



Transverse resolution for cosmic rays Ar+10%CO2 (B=0)



Compared to conventional readout, charge dispersion gives better resolution for the GEM and the Micromegas.

Charge dispersion TPC beam test in a magnet at KEK - October 2005

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	V.Lepeltier & Th.Zerguerras	LAL Orsay
<u>Germany</u>	R. Settles	MPI (Munich)
<u>Japan</u>	H.Kuroiwa & T.Takahashi	Hiroshima University
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	M.Habu, S.Matsushita, K.Nakamura & O.Nito	<u>Tokyo University of</u>
		<u>Agriculture & Technology</u>

KEK PS π2 beam test of Carleton & MT3 TPCs Beam data taken both in & outside the magnet for the two TPCs



•4 GeV/c hadrons (mostly πs^{i}) •0.5 & 1 GeV/c electrons •Super conducting 1.2 T magnet without return yoke Inner diameter : 850 mm •Effective length: 1 m



Carleton TPC in the beam outside the magnet

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The two beam test TPCs



- Micromegas 10 x10 cm²
- Drift distance: 16 cm
- 126 pads, 2 x 6 mm² each in 7 rows

-ALEPH preamps + 200 MHz FADCs rebinned to 25 MHz equivalent FADCs

- -Micromegas & GEMs 10 x10 cm²
- -Drift distance 25.9 cm
- 384 pads 2.3 x 6.3 mm² each in 16 rows

-ALEPH preamps + 11 MHz Aleph Time Projection Digitizers

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Track display - Ar+5%iC4H10 $Z_{drift} = 15.3 \text{ cm}$ Micromegas 2 x 6 mm² pads B = 1 TZ_{drift} = 15.3 \text{ cm}



Pad Response Function / Ar+5%iC4H10 Micromegas+Carleton TPC 2 x 6 mm² pads, B = 1 T

30 z regions / 0.5 cm step



Pad Response Function / Ar+5%iC4H10



Bias for central rows / Ar+5%iC4H10 B = 1 T



Transverse spatial resolution Ar+5%iC4H10 E=70V/cm D_{Tr} = 125 µm/√cm (Magboltz) @ B= 1T Micromegas+Carleton TPC 2 x 6 mm² pads



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Transverse resolution with no magnet - Angle dependence Ar+10% CO2, $D_{Tr} = 222 \mu/\sqrt{cm}$ (Magboltz) E=300 V/cm



MT3 TPC event display + Micromegas read out with Aleph TPDs 2.3 x 6.3 mm² pads Ar+5%iC4H10



Track display MT3 TPC with triple GEM readout

Part of MT3-TPC read out with Carleton FADCs 4 GeV/c π + beam 2.3 mm pitch x 6.3 mm pads 25 cm maximum drift distance Ar/CH4 (95/5) E = 50 V/cm D_{Tr} = 102 μ / \sqrt{cm} (Magboltz) @ 1T



Transverse resolution - MT3-TPC with Triple GEM Ar+5%CH4 E = 50 V/cm D_{Tr} = 102 µ/√cm (Magboltz) @ 1T



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Summary

- Traditional MPGD-TPC readout has difficulty achieving good resolution if wide pads similar in width to conventional wire TPC pads are used.
- With charge dispersion, the cluster charge can be dispersed in a controlled way such that wide pads can be used without sacrificing resolution. With such a readout system, we have achieved excellent resolution with wide pads both for the GEM and the Micromegas.
- The ILC-TPC resolution goal, 100 μm for all tracks, should be feasible with 2 x 6 mm² pads with a good understanding of systematics.
- R&D plans cosmic ray TPC tests at 4 T & two track resolution studies in a beam.
- R&D issues: New technology issues of fabrication & quality control, understanding/reducing bias. As charge dispersion pulses are slow, ~25 MHz digitizers could be used.



Conventional TPCs never achieve their potential! Example:Systematic effects in Aleph TPC at LEP



•Average Aleph resolution ~ 150 μ m •About 100 μ m best for all drift distances • Limit from diffusion σ (10 cm drift) ~ 20 μ m; σ (2 m drift) ~ 90 μ m